



**fcfm**

Ingeniería Eléctrica  
FACULTAD DE CIENCIAS  
FÍSICAS Y MATEMÁTICAS  
UNIVERSIDAD DE CHILE



# FI 2A2 ELECTROMAGNETISMO

## Clase 25

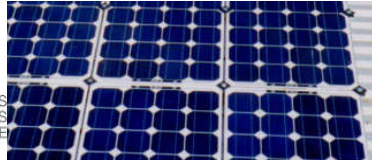
### Campos Variables en el Tiempo-IV

**LUIS S. VARGAS**  
Area de Energía  
Departamento de Ingeniería Eléctrica  
Universidad de Chile



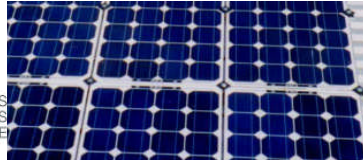
**fcfm**

Ingeniería Eléctrica  
FACULTAD DE CIENCIAS  
FÍSICAS Y MATEMÁTICAS  
UNIVERSIDAD DE CHILE



# INDICE

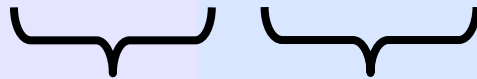
- Movimiento de cargas en presencia de campos
- Energía del campo eléctrico
- Energía del campo magnético
- Energía Electromagnética
- Fuerza
- Ejemplos



# Fuerza de Lorentz

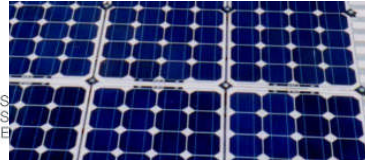
$$\vec{F} = q\vec{E} + q\vec{u} \times \vec{B}$$

Fuerza de Lorentz



Producida  
por campo  
eléctrico

Producida  
por campo  
magnético



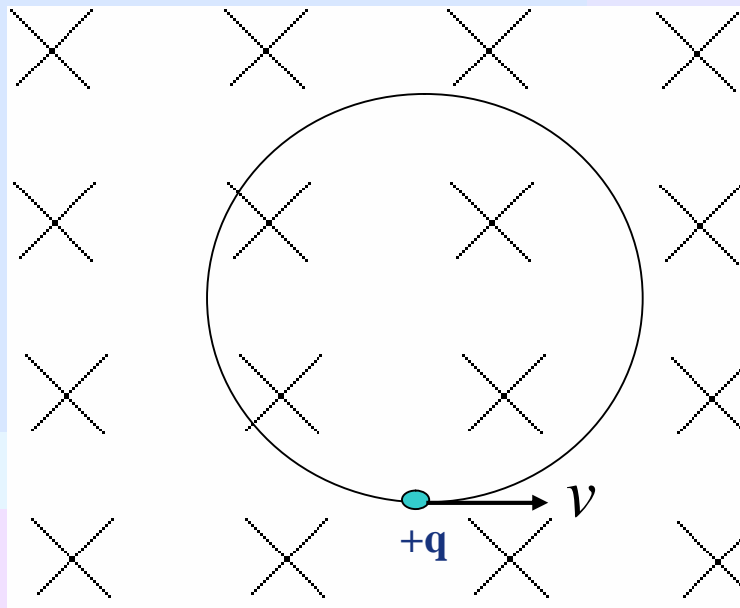
# Cargas en campos magnéticos

$$\vec{F} = q\vec{u} \times \vec{B}$$

Fuerza de Lorentz

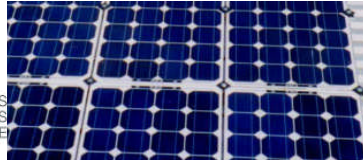
$$qvB = \frac{mv^2}{r}$$

$$\omega = \frac{v}{r} = \frac{qB}{m}$$



$\vec{B}$

Trayectoria circular

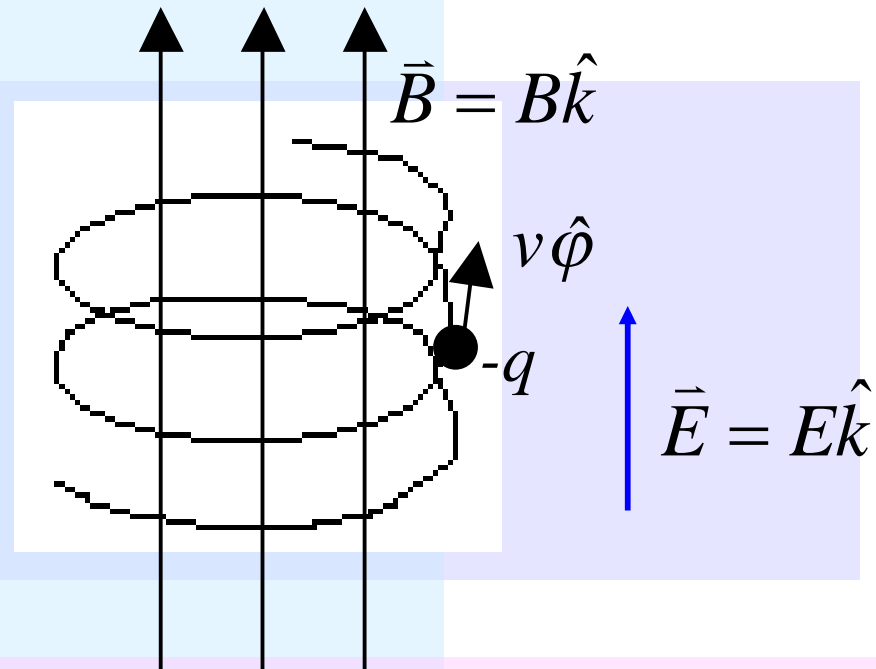


## Cargas en Campos Magnéticos

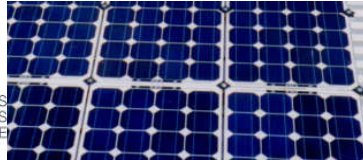
$$\vec{F} = q(\vec{E} + \vec{u} \times \vec{B})$$

$$qvB = \frac{mv^2}{r}$$

$$r = \frac{mv}{qB}$$



*Trayectoria helicoidal*

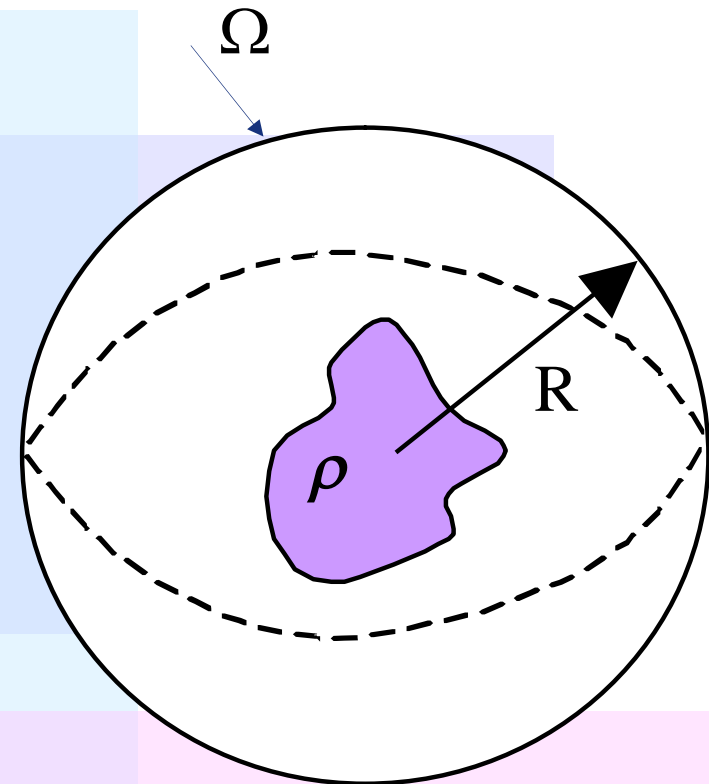


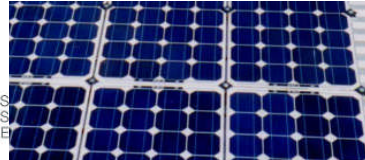
# Energía Electromagnética

Habíamos obtenido que la energía del campo eléctrico es

$$U = \frac{1}{2} \iiint_{\Omega} \vec{D} \cdot \vec{E} dv$$

Aquí  $\Omega$  es todo el espacio





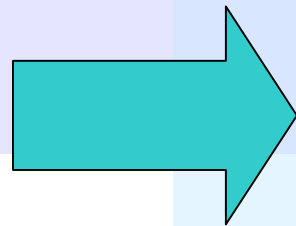
# Energía Electromagnética

**3ª Ecuación de Maxwell**

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad / \vec{H}.$$

**4ª Ecuación de Maxwell**

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \quad / \vec{E}.$$

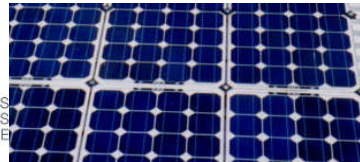


$$\vec{E} \cdot (\nabla \times \vec{H}) = \vec{E} \cdot \vec{J} + \vec{E} \cdot \frac{\partial \vec{D}}{\partial t}$$

$$\vec{H} \cdot (\nabla \times \vec{E}) = -\vec{H} \cdot \frac{\partial \vec{B}}{\partial t}$$

**Sumando**

$$\vec{E} \cdot (\nabla \times \vec{H}) - \vec{H} \cdot (\nabla \times \vec{E}) = \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} + \vec{H} \cdot \frac{\partial \vec{B}}{\partial t} + \vec{E} \cdot \vec{J}$$



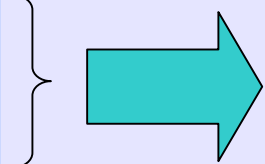
# Energía Electromagnética

$$\vec{E} \cdot (\nabla \times \vec{H}) - \vec{H} \cdot (\nabla \times \vec{E}) = \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} + \vec{H} \cdot \frac{\partial \vec{B}}{\partial t} + \vec{E} \cdot \vec{J}$$

Usando la propiedad  $\nabla \cdot (\vec{E} \times \vec{H}) = (\nabla \times \vec{E}) \cdot \vec{H} - (\nabla \times \vec{H}) \cdot \vec{E}$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

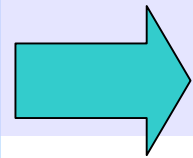
$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$



$$\nabla \cdot (\vec{E} \times \vec{H}) = -\vec{E} \cdot \left( \frac{\partial \vec{D}}{\partial t} \right) - \vec{H} \cdot \left( \frac{\partial \vec{B}}{\partial t} \right) - \vec{E} \cdot \vec{J}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{B} = \mu \vec{H}$$

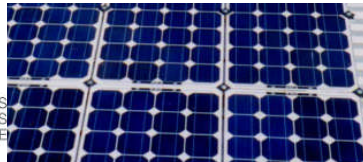


$$\nabla \cdot (\vec{E} \times \vec{H}) = -\vec{E} \cdot \left( \frac{\partial (\epsilon \vec{E})}{\partial t} \right) - \mu \vec{H} \cdot \left( \frac{\partial \vec{B}}{\partial t} \right) - \vec{E} \cdot \vec{J}$$

$$\nabla \cdot (\vec{E} \times \vec{H}) = -\epsilon \frac{\partial}{\partial t} (\vec{E} \cdot \vec{E}) - \mu \frac{\partial}{\partial t} (\vec{B} \cdot \vec{B}) - \vec{E} \cdot \vec{J}$$

$$\nabla \cdot (\vec{E} \times \vec{H}) = -\frac{\partial}{\partial t} (\vec{E} \cdot \epsilon \vec{E}) - \frac{\partial}{\partial t} (\mu \vec{B} \cdot \vec{B}) - \vec{E} \cdot \vec{J}$$





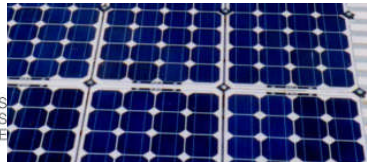
## Energía Electromagnética

$$\Rightarrow -\nabla \cdot (\vec{E} \times \vec{H}) = \frac{1}{2} \cdot \frac{\partial(\vec{E} \cdot \vec{D})}{\partial t} + \frac{1}{2} \cdot \frac{\partial(\vec{H} \cdot \vec{B})}{\partial t} + \vec{E} \cdot \vec{J}$$

$$\Rightarrow \frac{1}{2} \cdot \frac{\partial(\vec{E} \cdot \vec{D})}{\partial t} + \frac{1}{2} \cdot \frac{\partial(\vec{H} \cdot \vec{B})}{\partial t} + \vec{E} \cdot \vec{J} = -\nabla \cdot (\vec{E} \times \vec{H})$$

Tomando la integral sobre un volumen  $\Omega$  muy grande

$$\Rightarrow \iiint_{\Omega} \frac{1}{2} \frac{\partial}{\partial t} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv + \iiint_{\Omega} \vec{E} \cdot \vec{J} dv = -\iiint_{\Omega} \nabla \cdot (\vec{E} \times \vec{H}) dv$$



# Energía Electromagnética

Nos interesa el caso cuando  $\Omega$  es todo el espacio

$$\Rightarrow \iiint_{\Omega} \frac{1}{2} \frac{\partial}{\partial t} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv + \iiint_{\Omega} \vec{E} \cdot \vec{J} dv = - \iiint_{\Omega} \underbrace{\nabla \cdot \vec{E}}_{\propto \frac{1}{r^2}} \cdot \underbrace{(\vec{E} \times \vec{H})}_{\propto r^2} dv$$

$$\propto \frac{1}{r^3} \quad \therefore \lim_{r \rightarrow \infty} (\iiint_{\Omega} \dots) = 0$$



# Energía Electromagnética

Nos interesa el caso cuando  $\Omega$  es todo el espacio

$$\Rightarrow \underbrace{\frac{\partial}{\partial t} \iiint_{\Omega} \frac{1}{2} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv}_{\text{Energía del campo electromagnético}} + \underbrace{\iiint_{\Omega} \vec{E} \cdot \vec{J} dv}_{\text{Potencia consumida por efecto Joule}} = 0$$

Energía del campo electromagnético

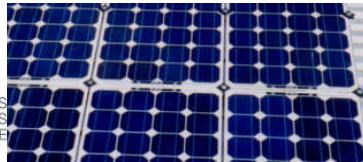
Potencia consumida por efecto Joule

$$U = \iiint_{\Omega} \frac{1}{2} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv$$

$$P_{\text{Joule}} = \iiint_{\Omega} \vec{E} \cdot \vec{J} dv$$

Energía del campo eléctrico

Energía del campo magnético



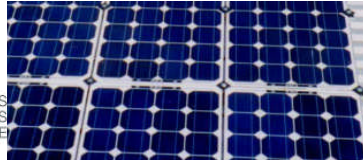
## Energía Electromagnética

Si no hay pérdidas Joule  $P_{Joule} = \iiint_{\Omega} \vec{E} \cdot \vec{J} dv = 0$

$$\Rightarrow \frac{\partial}{\partial t} \iiint_{\Omega} \frac{1}{2} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv = 0$$

$$\therefore U = \iiint_{\Omega} \frac{1}{2} (\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B}) dv = Cte.$$

**En ausencia de pérdidas Joule la energía electromagnética se conserva**



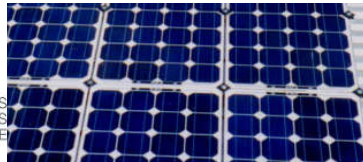
# Energía Electromagnética

Si no hay campos magnéticos

$$U = \iiint_{\Omega} \frac{1}{2} (\vec{E} \cdot \vec{D}) dv$$

Si no hay campos eléctricos

$$U = \iiint_{\Omega} \frac{1}{2} (\vec{H} \cdot \vec{B}) dv$$

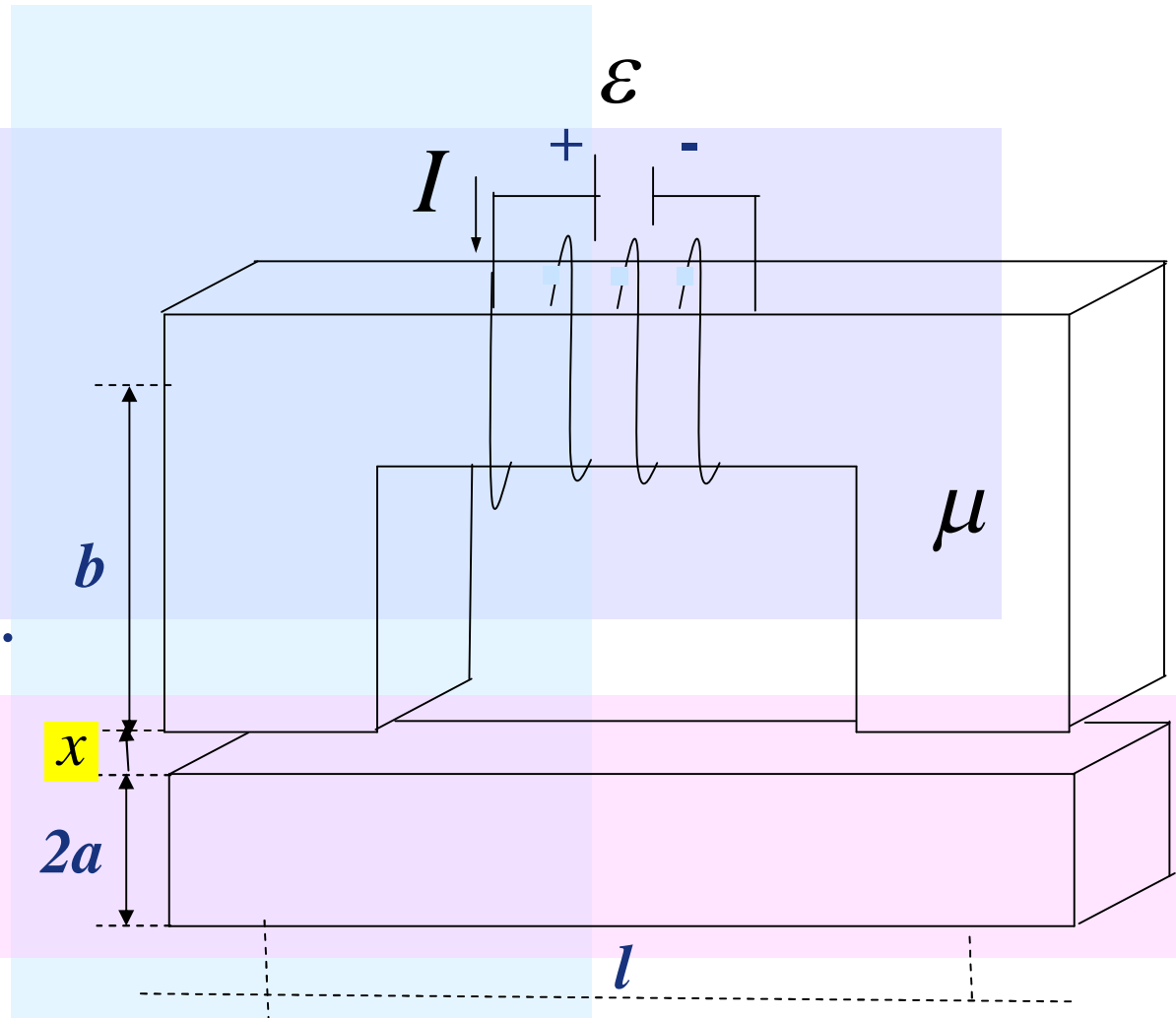


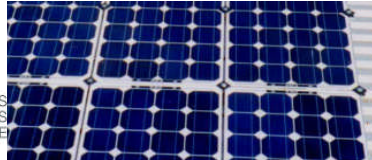
# Energía Electromagnética

## Ejemplo 1

Calcular la energía para el sistema compuesto de un material ferromagnético de sección cuadrada  $S$ .  
Aproximar para el caso

$$\mu \gg \mu_0$$





# Energía Electromagnética

La energía es

$$U = \iiint_{\Omega} \frac{1}{2} (\vec{H} \cdot \vec{B}) dv$$

Tomando la Ley  
Circuital de Ampere  
en el camino medio

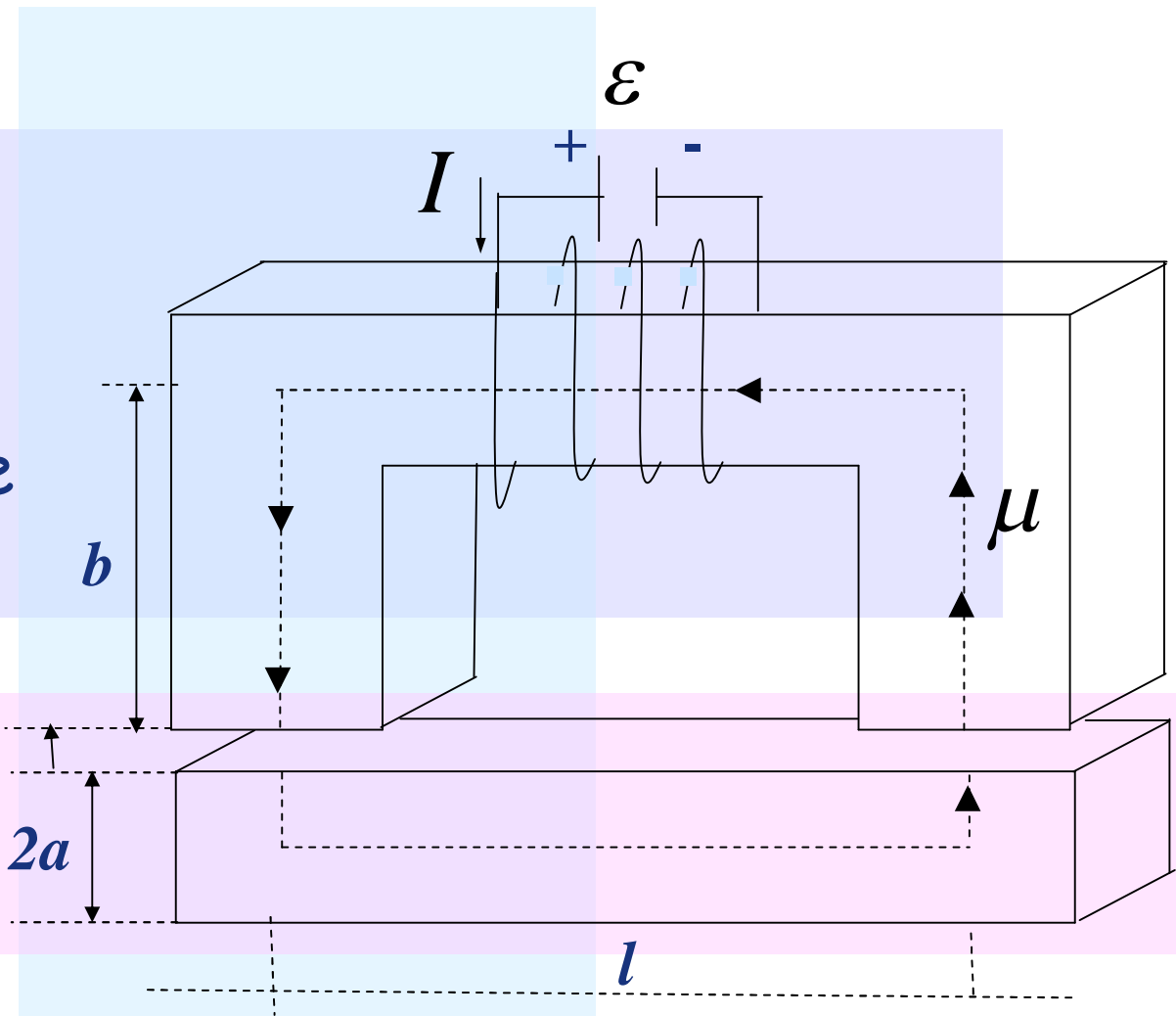
$$\oint \vec{H} \cdot d\vec{l} = I_{total}$$

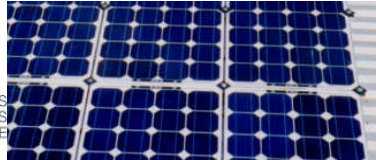
$$I_{total} = -NI$$

$x, \hat{i}$

$2a$

$l$

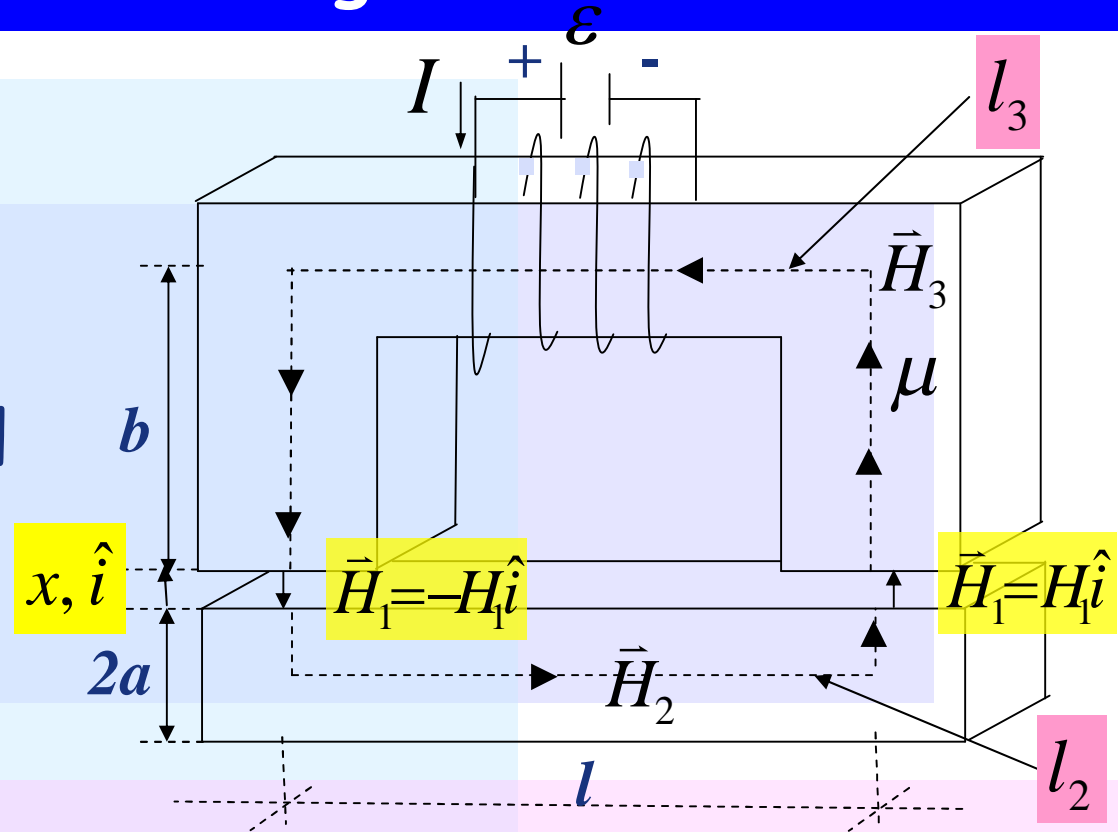




# Energía Electromagnética

Supondremos:

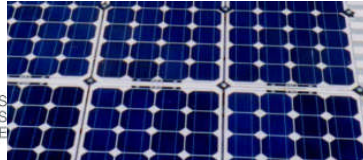
- B y H constantes en cada medio
- B y H son rectilíneos y paralelos al material (se desprecian efectos de borde)



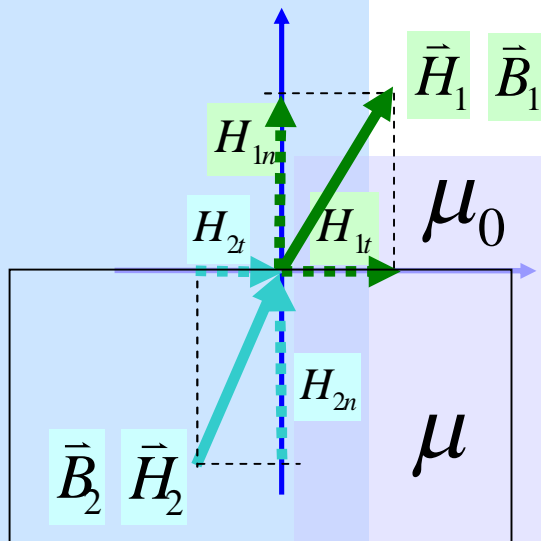
$$\oint \vec{H} \cdot d\vec{l} = \int \vec{H}_1 \cdot dx \hat{i} + \int_{l_3} \vec{H}_3 \cdot d\vec{l} + \int (-H_1) \cdot dx (-\hat{i}) + \int_{l_2} \vec{H}_2 \cdot d\vec{l}$$

$$\oint \vec{H} \cdot d\vec{l} = H_1 x + H_3 l_3 + H_1 x + H_2 l_2$$





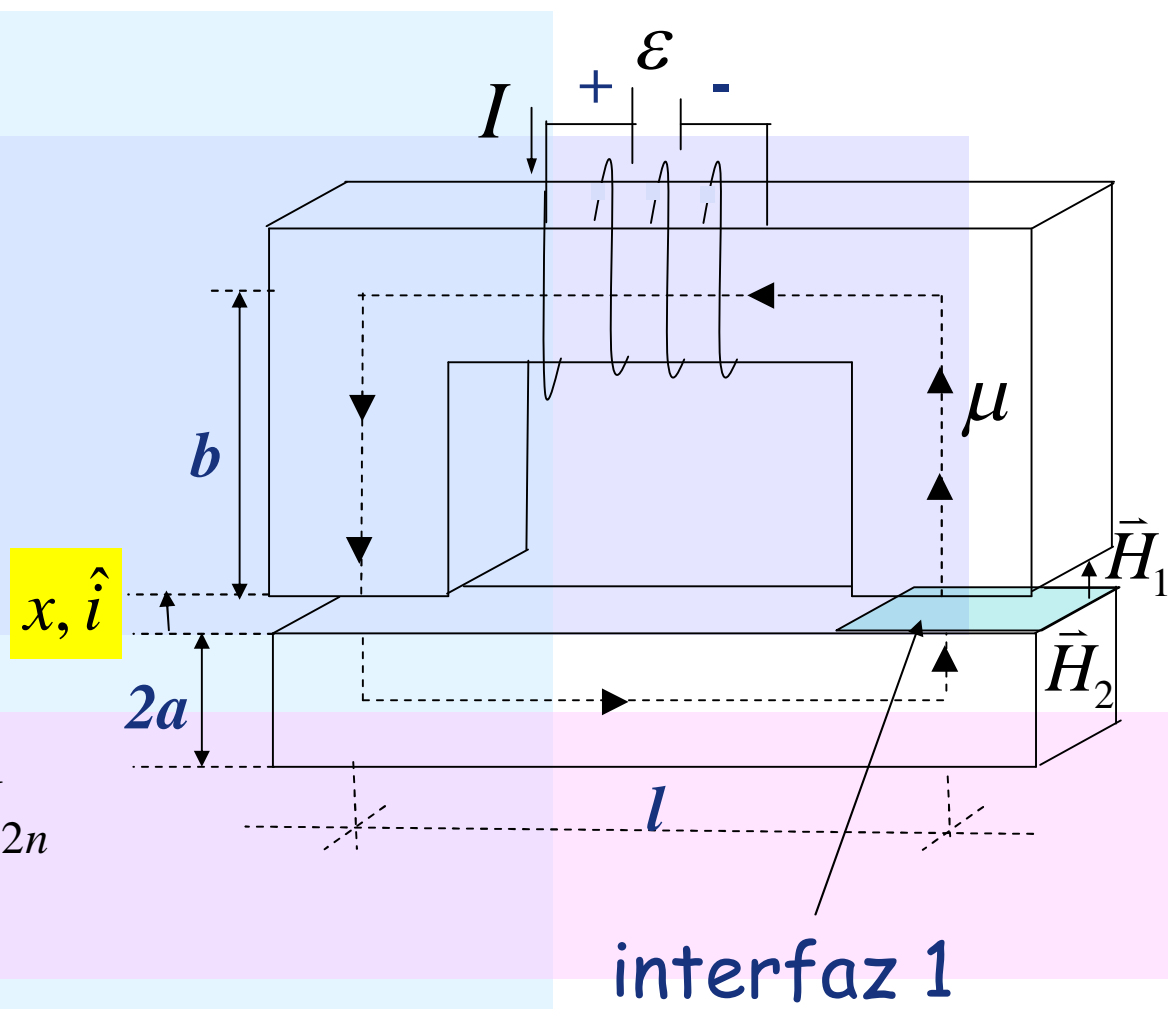
# Energía Electromagnética

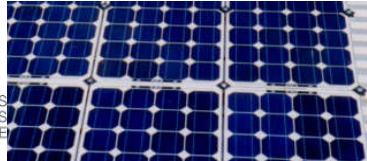


En la interfaz 1

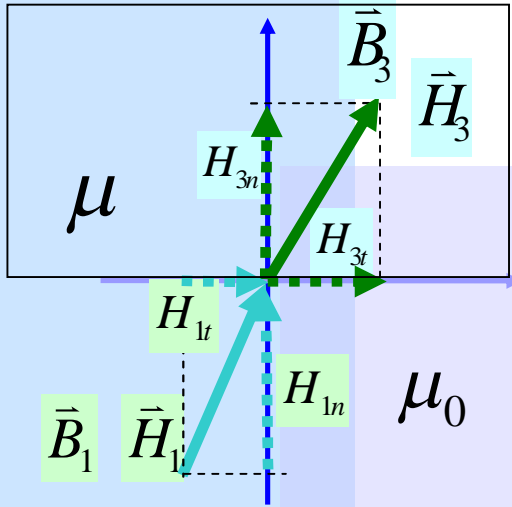
$$B_{1n} = B_{2n} \Leftrightarrow \mu_0 H_{1n} = \mu H_{2n}$$

$$\therefore \mu_0 H_1 = \mu H_2$$





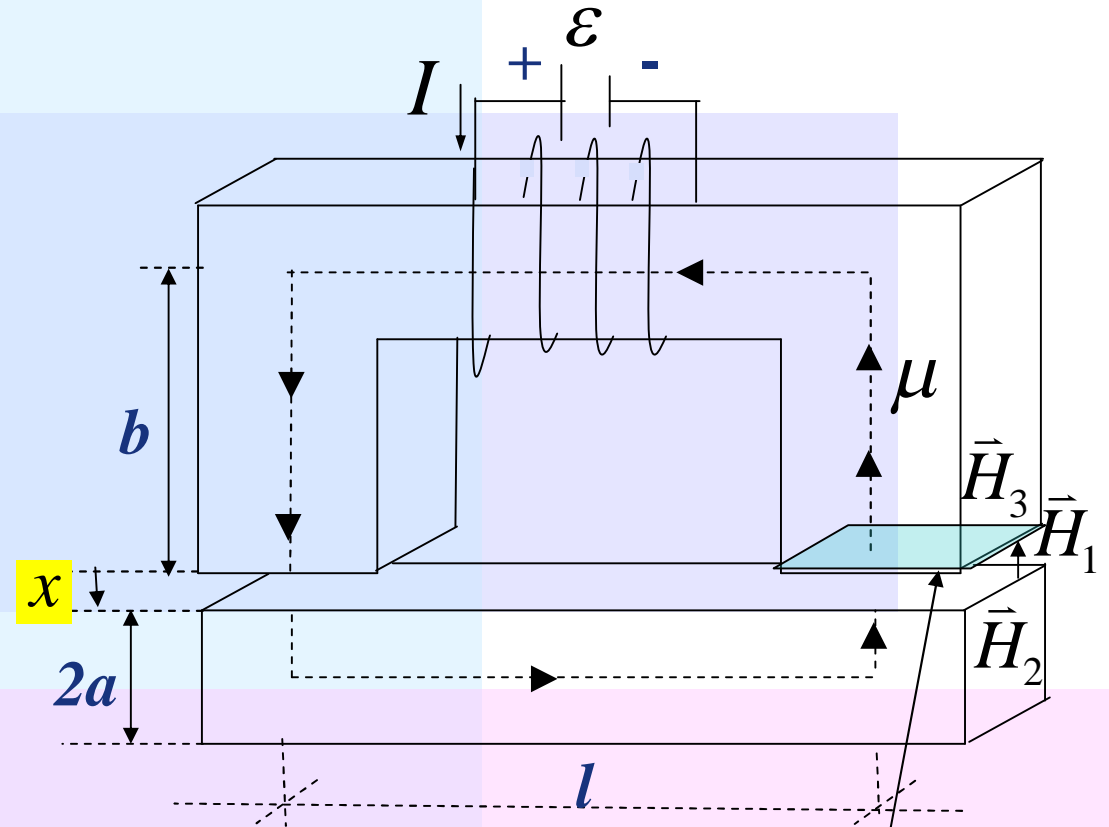
# Energía Electromagnética



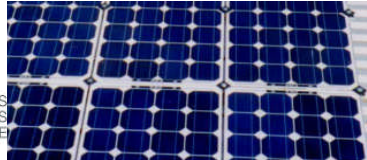
En la interfaz 2  
se cumple

$$B_{1n} = B_{3n} \Leftrightarrow \mu_0 H_{1n} = \mu H_{3n}$$

$$\therefore \mu_0 H_1 = \mu H_3 \quad \Rightarrow \quad \vec{H}_2 = \vec{H}_3$$



interfaz 2



# Energía Electromagnética

$$\oint \vec{H} \cdot d\vec{l} = H_1 x + H_3 l_3 + H_1 x + H_2 l_2$$

$$\oint \vec{H} \cdot d\vec{l} = 2H_1 x + H_2 (l_2 + l_3)$$

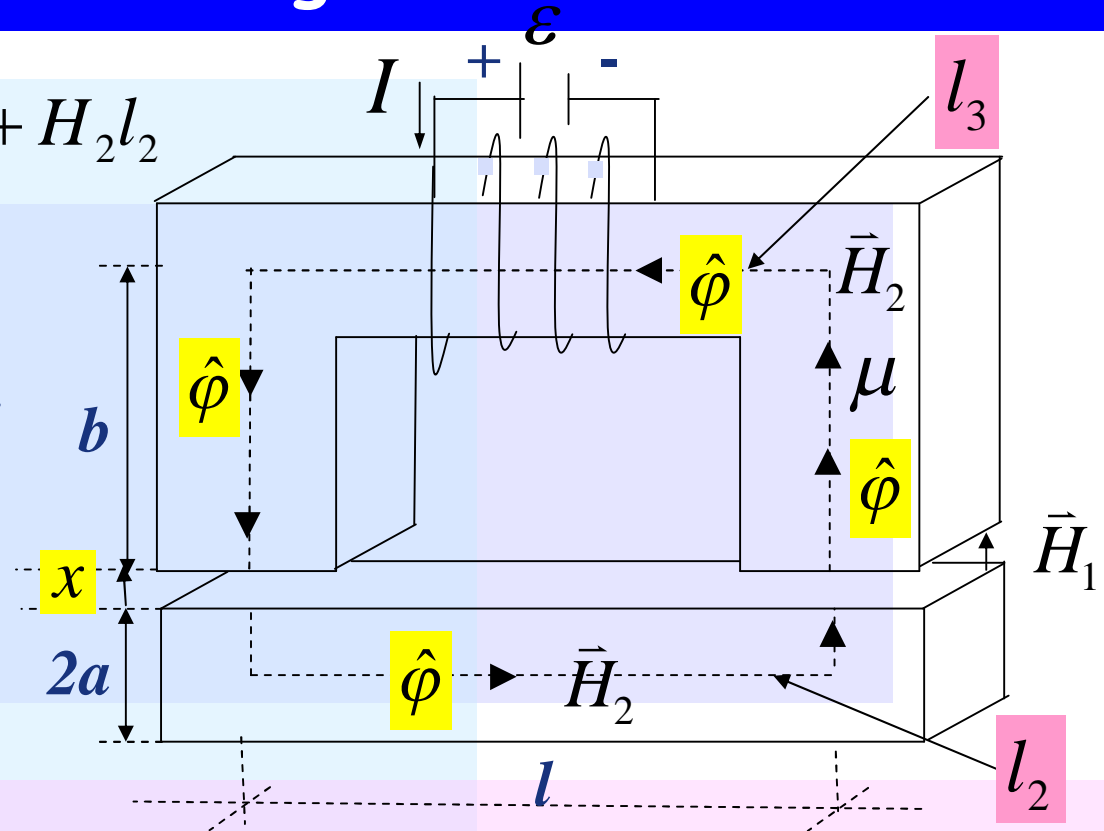
Reemplazando en la ley  
circuital de Ampere

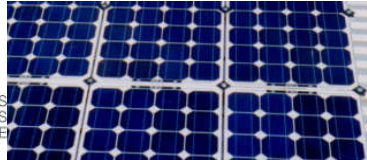
$$2H_1 x + H_2 (l_2 + l_3) = -NI$$

Teniamos  $\mu_0 H_1 = \mu H_2$

$$\vec{H}_1 = -\frac{\mu NI}{2\mu x + \mu_0 (l_2 + l_3)} \hat{\phi}$$

$$\vec{H}_2 = -\frac{\mu_0 NI}{2\mu x + \mu_0 (l_2 + l_3)} \hat{\phi}$$





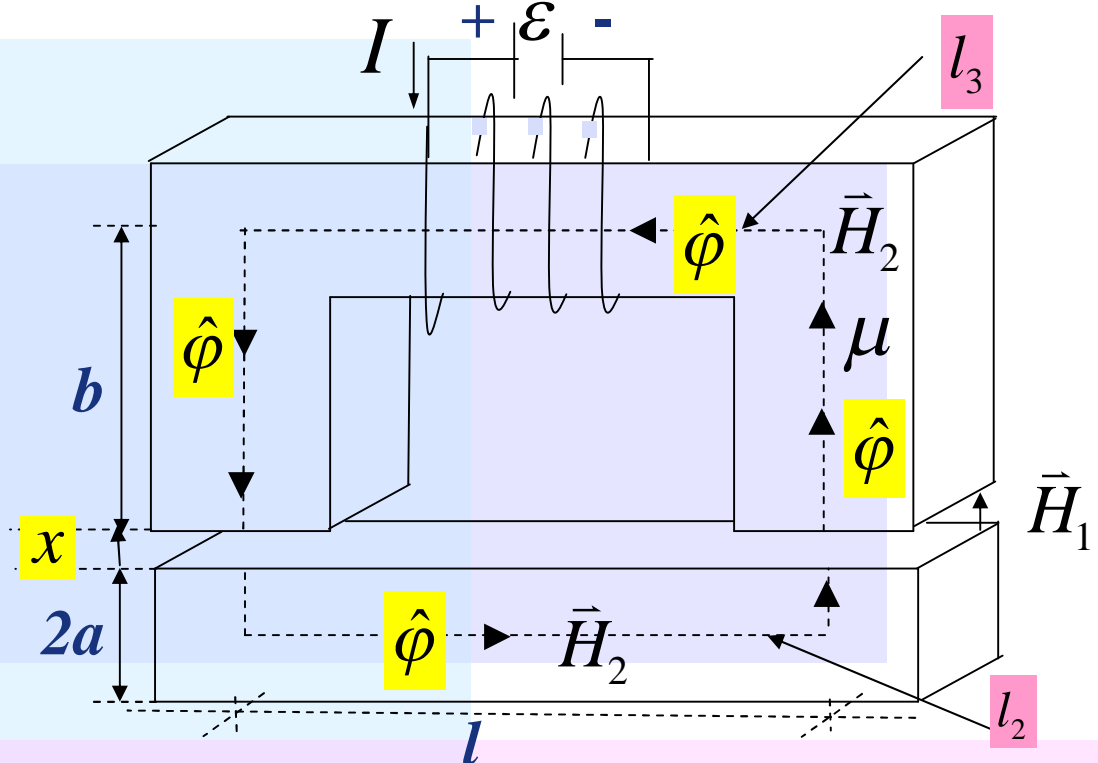
# Energía Electromagnética

$$\vec{H}_1 = -\frac{\mu NI}{2\mu x + \mu_0(l_2 + l_3)} \hat{\phi}$$

$$\vec{B}_1 = -\frac{\mu_0 \mu NI}{2\mu x + \mu_0(l_2 + l_3)} \hat{\phi}$$

$$\vec{H}_2 = -\frac{\mu_0 NI}{2\mu x + \mu_0(l_2 + l_3)} \hat{\phi}$$

$$\vec{B}_2 = -\frac{\mu \mu_0 NI}{2\mu x + \mu_0(l_2 + l_3)} \hat{\phi}$$



$$U = \iiint_{\Omega} \frac{1}{2} (\vec{H} \cdot \vec{B}) dv \Rightarrow U = \frac{1}{2} H_1 B_1 Vol_1 + \frac{1}{2} H_2 B_2 Vol_2$$

$$Vol_1 = 2xS$$
$$Vol_2 = S(l_2 + l_3)$$

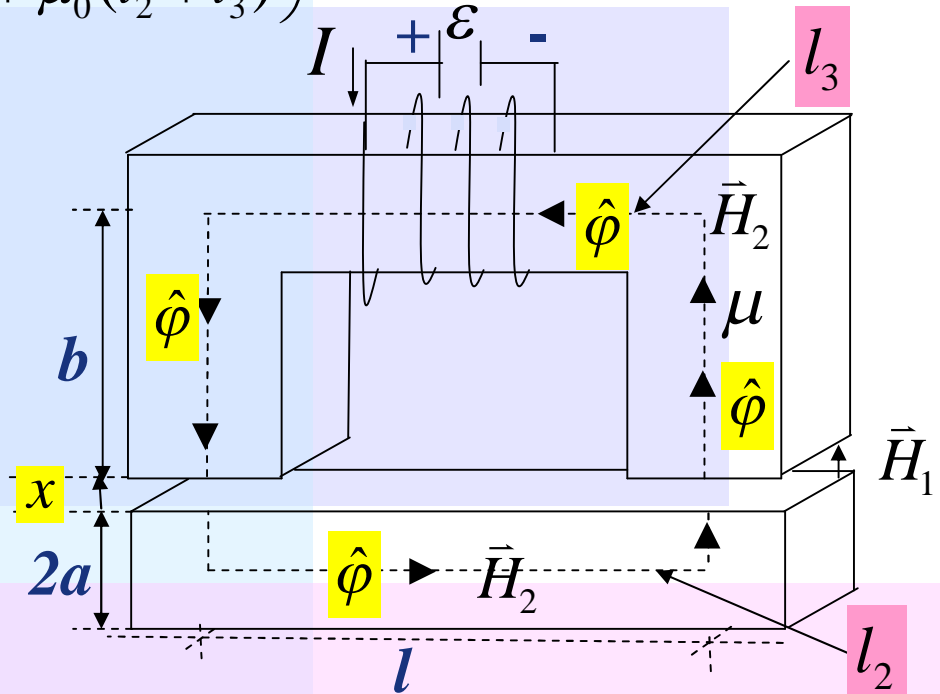
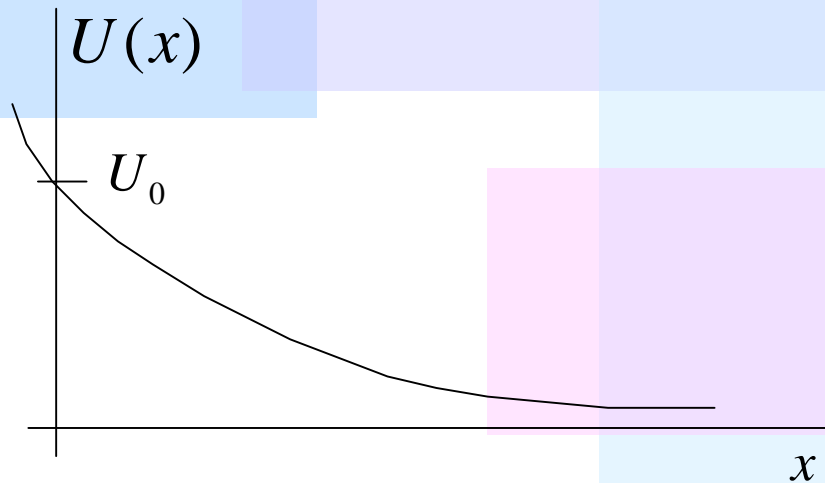
$$\Rightarrow U = \frac{1}{2} \mu_0 \mu^2 \left( \frac{NI}{2\mu x + \mu_0(l_2 + l_3)} \right)^2 2xS + \frac{1}{2} \mu_0^2 \mu \left( \frac{NI}{2\mu x + \mu_0(l_2 + l_3)} \right)^2 S(l_2 + l_3)$$

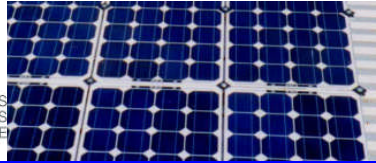


# Energía Electromagnética

$$\Rightarrow U = \frac{1}{2} S \mu_0 \mu (2x\mu + \mu_0 (l_2 + l_3)) \left( \frac{NI}{2\mu x + \mu_0 (l_2 + l_3)} \right)^2$$

$$\therefore U = \frac{1}{2} \left( \frac{S \mu_0 \mu N^2 I^2}{2\mu x + \mu_0 (l_2 + l_3)} \right)$$





# Energía Electromagnética

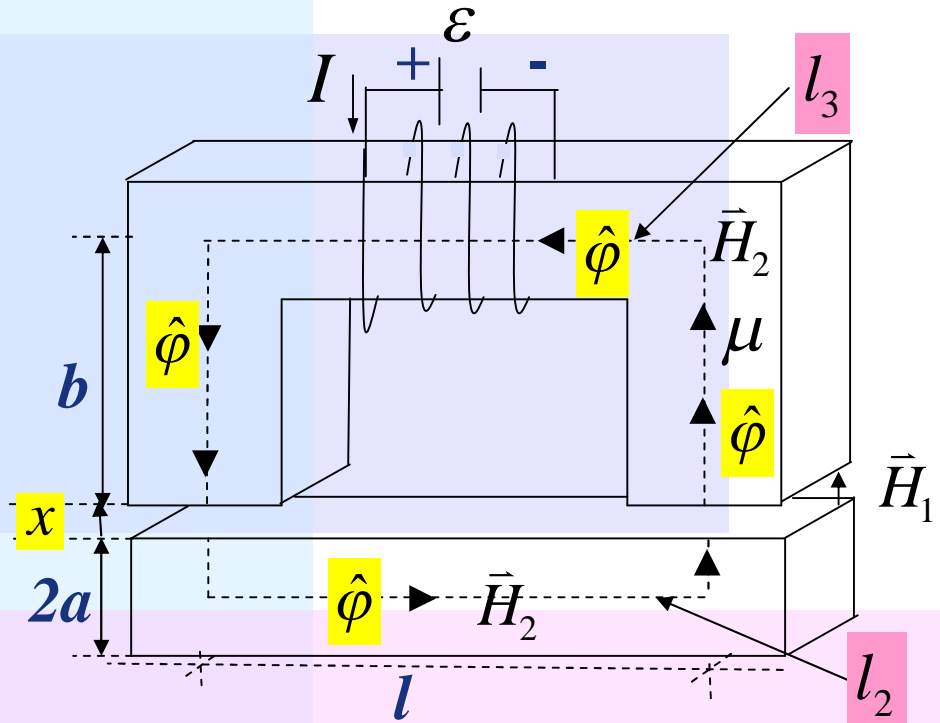
$$\text{Si } \mu \gg \mu_0 \Rightarrow \mu_R \mu_0 \gg \mu_0$$

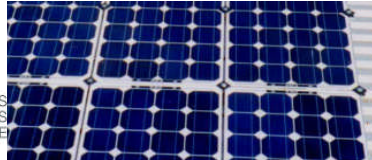
$$U = \frac{1}{2} H_1 B_1 Vol_1 + \frac{1}{2} H_2 B_2 Vol_2$$

$$U = \frac{1}{2\mu_0} B_1^2 Vol_1 + \frac{1}{2\mu} B_2^2 Vol_2$$

$$B_1 = B_2$$

$$\Rightarrow U = \frac{1}{2\mu_0} B_1^2 \left( Vol_1 + \frac{1}{2\mu_R} Vol_2 \right)$$





# Energía Electromagnética

$$U = \frac{1}{2\mu_0} B_1^2 \left( Vol_1 + \frac{1}{2\mu_R} Vol_2 \right)$$

Notar que  $Vol_1 = 2Sx$

$$Vol_2 \approx S(l_2 + l_3)$$

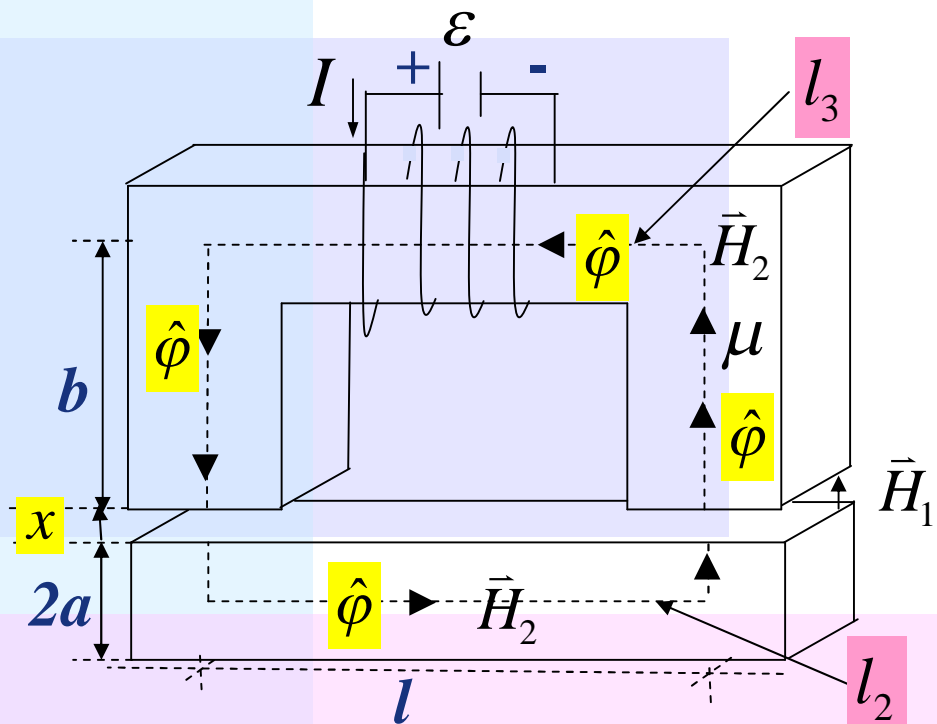
Típicamente

$$Vol_2 \approx 10 \times Vol_1$$

$$\mu_R > 1000$$

$$\Rightarrow U \cong \frac{1}{2\mu_0} B_1^2 Vol_1$$

La mayor parte de la energía se concentra en el entrehierro





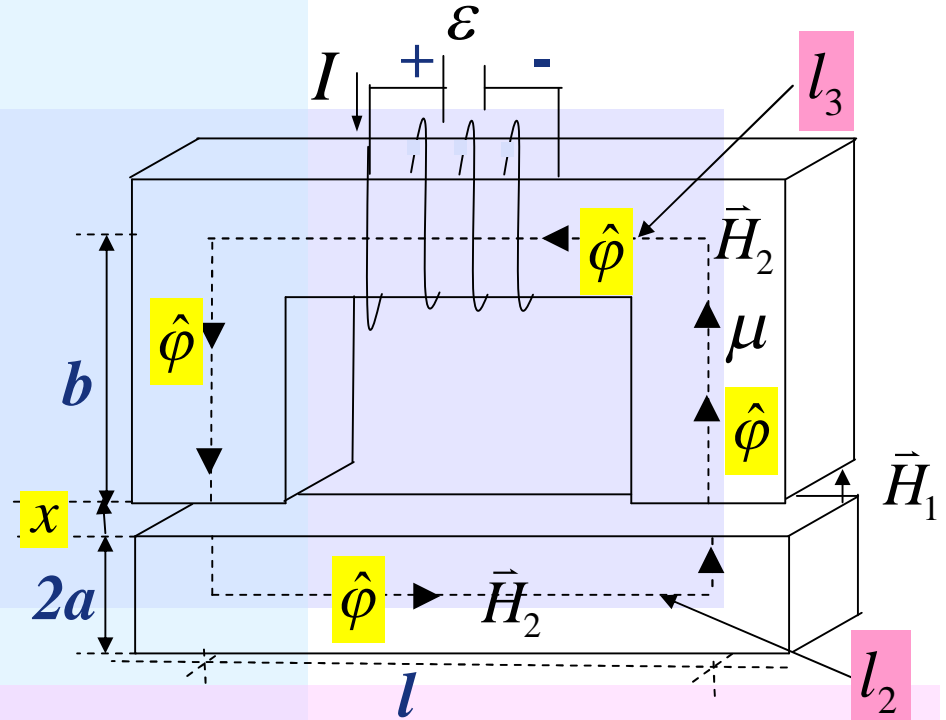
# Energía Electromagnética

$$\Rightarrow U \cong \frac{1}{2\mu_0} B_1^2 Vol_1$$

$$\therefore U \cong xS\mu_0 \left( \frac{\mu NI}{2\mu x + \mu_0(l_2 + l_3)} \right)^2$$

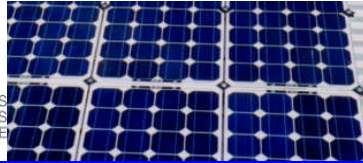
$$U \cong xS\mu_0 \left( \frac{NI}{2x + \frac{\mu_0}{\mu}(l_2 + l_3)} \right)^2$$

$$\Rightarrow U \cong xS\mu_0 \left( \frac{NI}{2x} \right)^2 = S\mu_0 \frac{N^2 I^2}{4x}$$



Fórmula aproximada





# Fuerza Electromagnética

Trabajo es igual al cambio de energía

$$-\vec{F} \cdot d\vec{l} = dU$$

$$\vec{F} = -\frac{dU}{dl} \hat{l}$$

En general

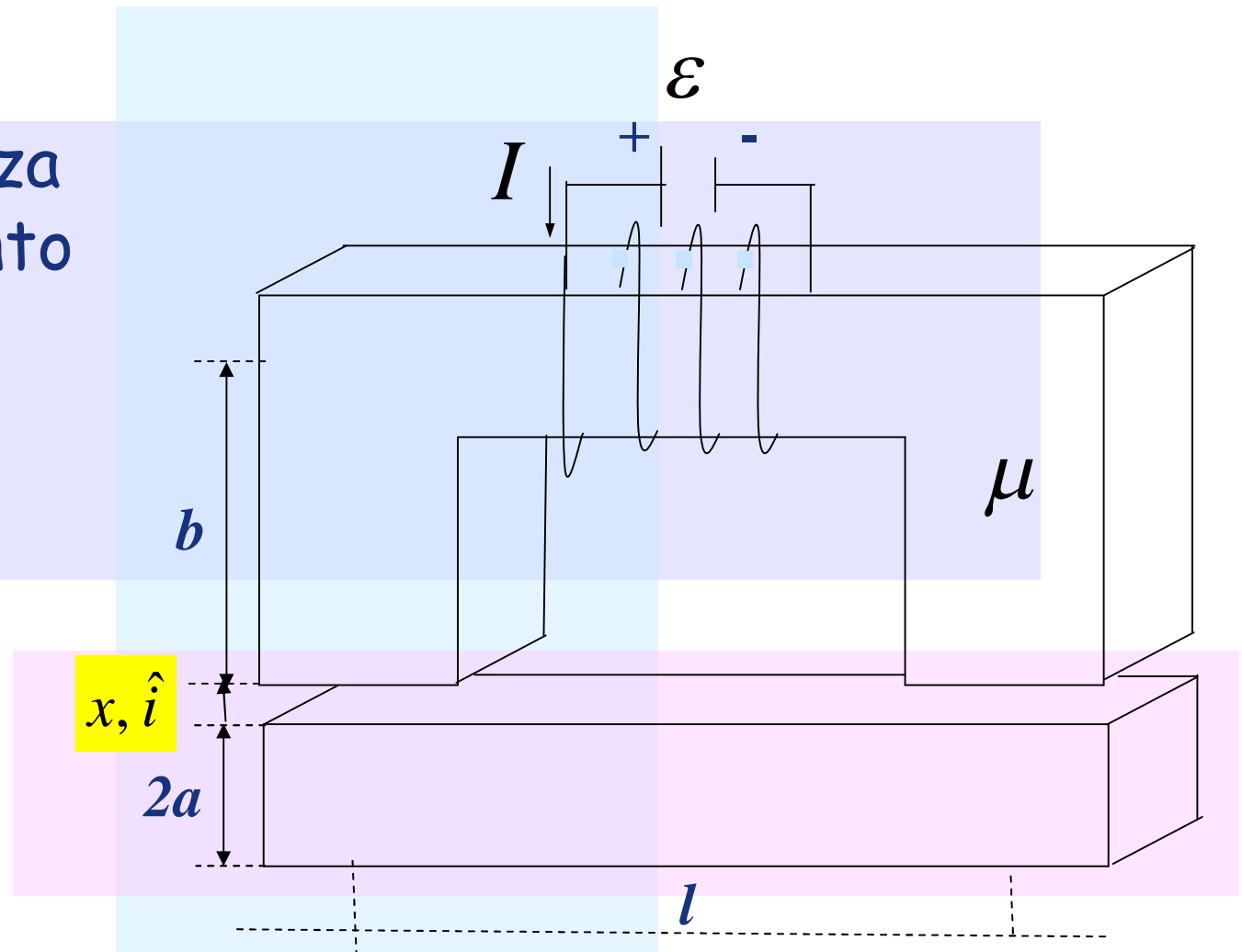
$$\vec{F} = -\nabla U$$



# Fuerza Electromagnética

## Ejemplo 2

Calcular la Fuerza sobre el segmento inferior





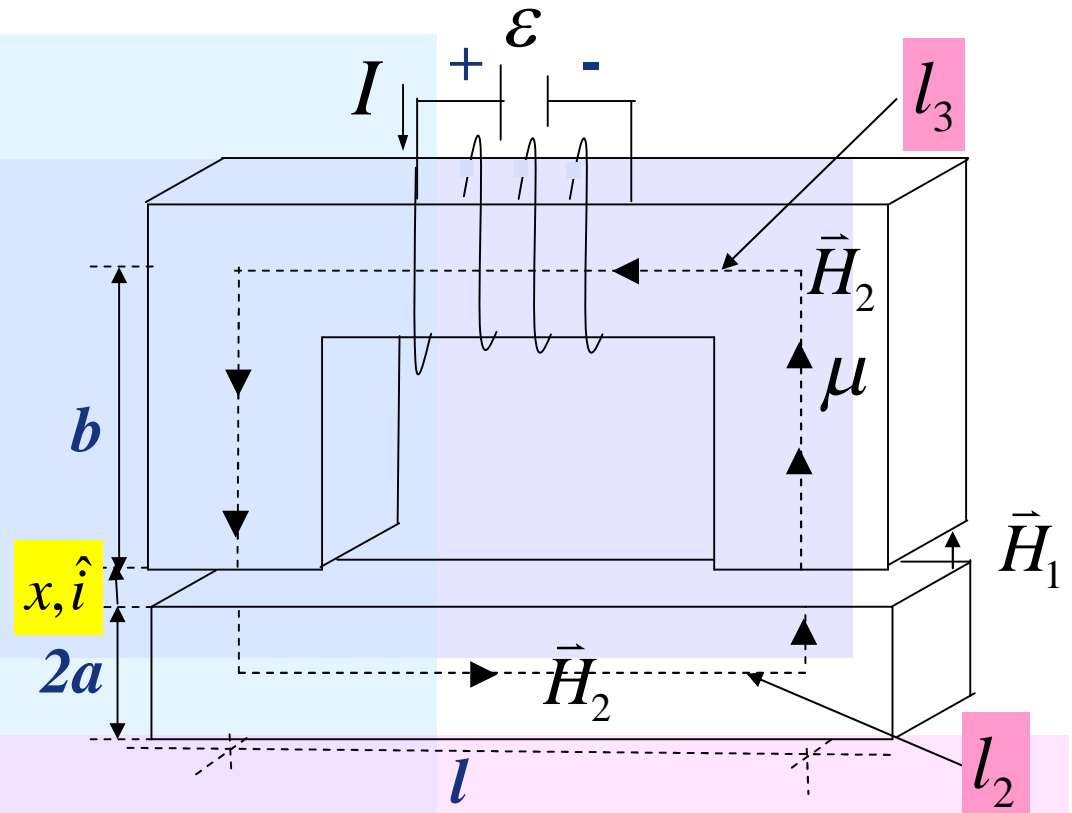
# Fuerza Electromagnética

$$\vec{F} = -\nabla U$$

$$\vec{F} = -\frac{\partial}{\partial x} \frac{1}{2} \left( \frac{\mu\mu_0 N^2 S}{2\mu x + \mu_0(l_2 + l_3)} \right) I^2 \hat{i}$$

$$\vec{F} = -\frac{1}{2} \left( -\frac{\mu\mu_0 N^2 S I^2}{(2\mu x + \mu_0(l_2 + l_3))^2} \right) 2\mu \hat{i}$$

$$\vec{F} = \left( \frac{\mu_0 \mu^2 N^2 S I^2}{[2\mu x + \mu_0(l_2 + l_3)]^2} \right) \hat{i}$$



Fuerza tiende a acercar la barra inferior: Electroimán.  
Notar que hemos supuesto corriente constante.



# Fuerza Electromagnética

Usando la fórmula aproximada

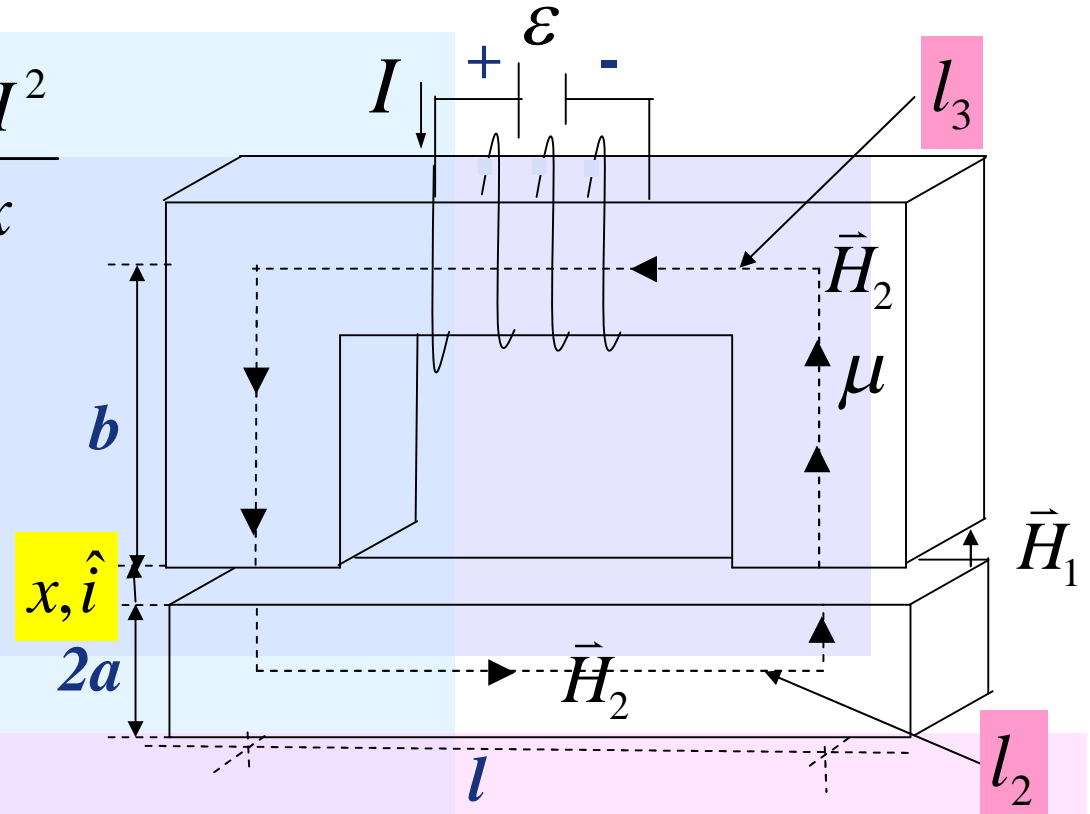
$$U = S\mu_0 \frac{N^2 I^2}{4x}$$

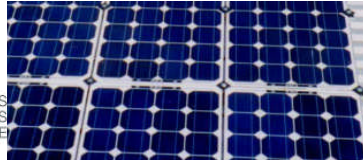
$$\vec{F} = -\nabla U$$

$$\vec{F} = -\frac{\partial}{\partial x} \left[ S\mu_0 \frac{N^2 I^2}{4x} \right] \hat{i}$$

Asumiendo que la corriente es constante

$$\vec{F} = S\mu_0 \frac{N^2 I^2}{4x^2} \hat{i}$$





# Energía de Campo Magnético

Usando la ecuación de la potencia

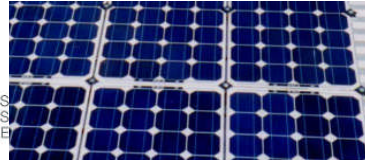
$$U = \int_{t_0}^t P(t) dt = \int_{t_0}^t v(t) i(t) dt$$

$$v(t) = \frac{\partial \Phi}{\partial t} \Rightarrow d\Phi = v(t) dt \quad \Rightarrow U = \int_{\phi_0}^{\phi} i(t) d\Phi$$

como  $\Phi = Li \quad \Rightarrow U = \int_{\phi_0}^{\phi} \frac{\Phi}{L} d\Phi = \frac{1}{2} \frac{\Phi^2}{L}$

Se obtiene

$$U = \frac{1}{2} \frac{\Phi^2}{L} = \frac{1}{2} Li^2$$

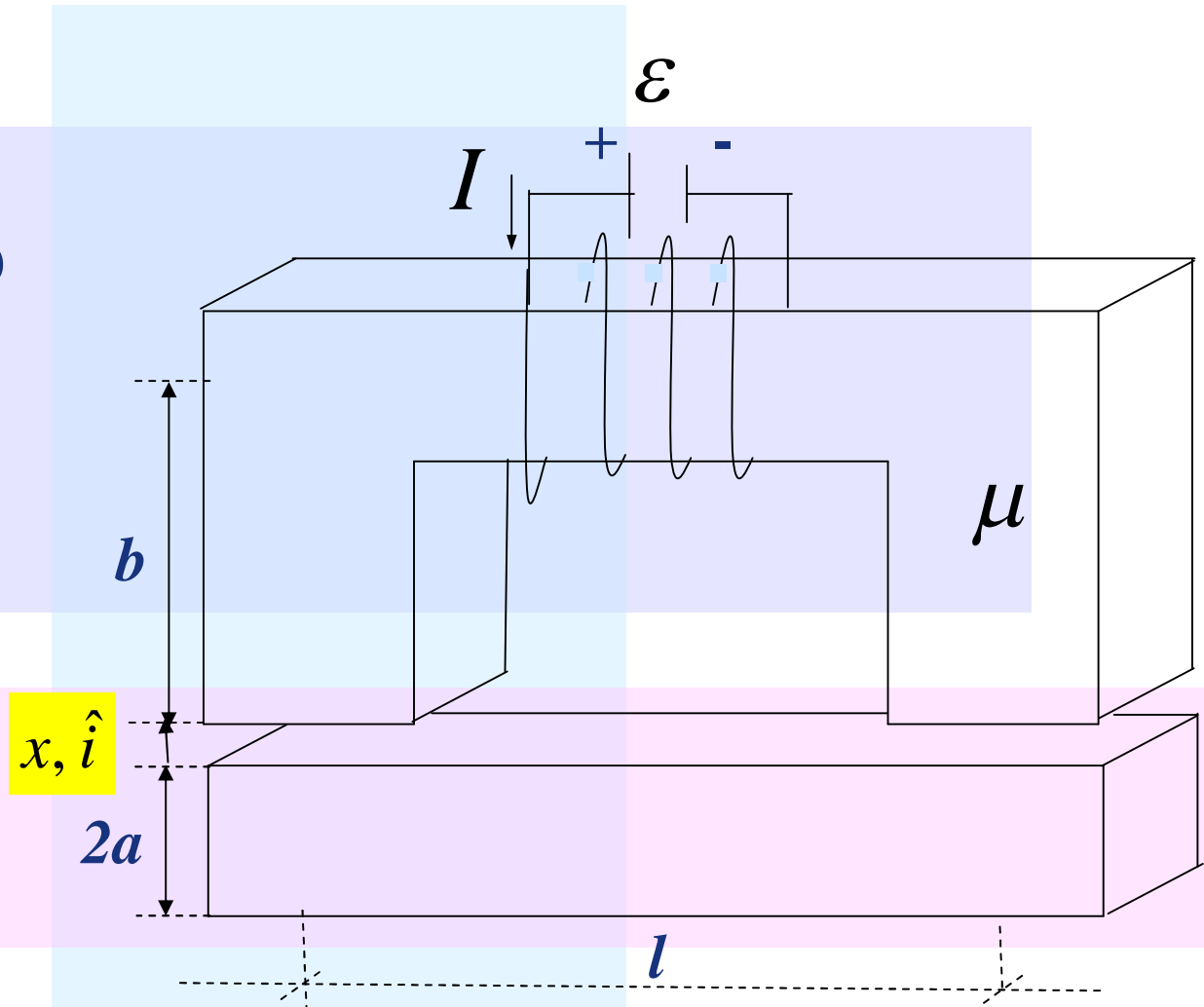


# Energía de Campo Magnético

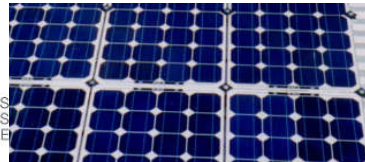
## Ejemplo 3

Calcular la Fuerza sobre el segmento inferior usando la ecuación de energía

$$U = \frac{1}{2} \frac{\Phi^2}{L} = \frac{1}{2} Li^2$$







# Energía de Campo Magnético

$$U = \frac{1}{2} LI^2$$

$$\Rightarrow U = \frac{1}{2} \left( \frac{\mu\mu_0 N^2 S}{2\mu x + \mu_0(l_2 + l_3)} \right) I^2$$

$$\vec{F} = -\frac{\partial}{\partial x} \frac{1}{2} \left( \frac{\mu\mu_0 N^2 S}{2\mu x + \mu_0(l_2 + l_3)} \right) I^2 \hat{i}$$

$$\vec{F} = -\frac{1}{2} \left( -\frac{\mu\mu_0 N^2 S I^2}{(2\mu x + \mu_0(l_2 + l_3))^2} \right) 2\mu \hat{i}$$

$$\vec{F} = \left( \frac{\mu_0 \mu^2 N^2 S I^2}{(2\mu x + \mu_0(l_2 + l_3))^2} \right) \hat{i}$$

